

# What regulates electron injection in diffusive shock acceleration?



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## Abstract

Collisionless shocks are one of the most efficient sources for energetic nonthermal particles. Although the microphysics of proton acceleration is revealed to some extent, the exact mechanisms that channel a fraction of thermal electrons to nonthermal population are still not fully understood.

The key open questions are

- What are the crucial processes and threshold condition that determine electron injection in diffusive shock acceleration (DSA)?
- How does the acceleration efficiency depend on the shock speed ( $v_{sh}$ ), Alfvén Mach number ( $M_A = v_{sh}/v_A$ ), and the sonic Mach number ( $M_s = v_{sh}/v_{th}$ )?
- What are the effects of upstream magnetic field inclination ( $\theta_{Bn}$ )?

To develop a comprehensive theory of electron acceleration, we have performed a survey of fully kinetic non-relativistic shock simulations in spatially 1D geometry using the massively parallel electromagnetic Particle-In-Cell code, Tristan-MP. The results are crucial to understand the nonthermal phenomenology of a variety of heliophysical and astrophysical collisionless shocks from interstellar space to galaxy clusters.

## Journey from thermal to nonthermal

Low Mach number shock

- Electrons escape from the shock.
- Acceleration stalls.

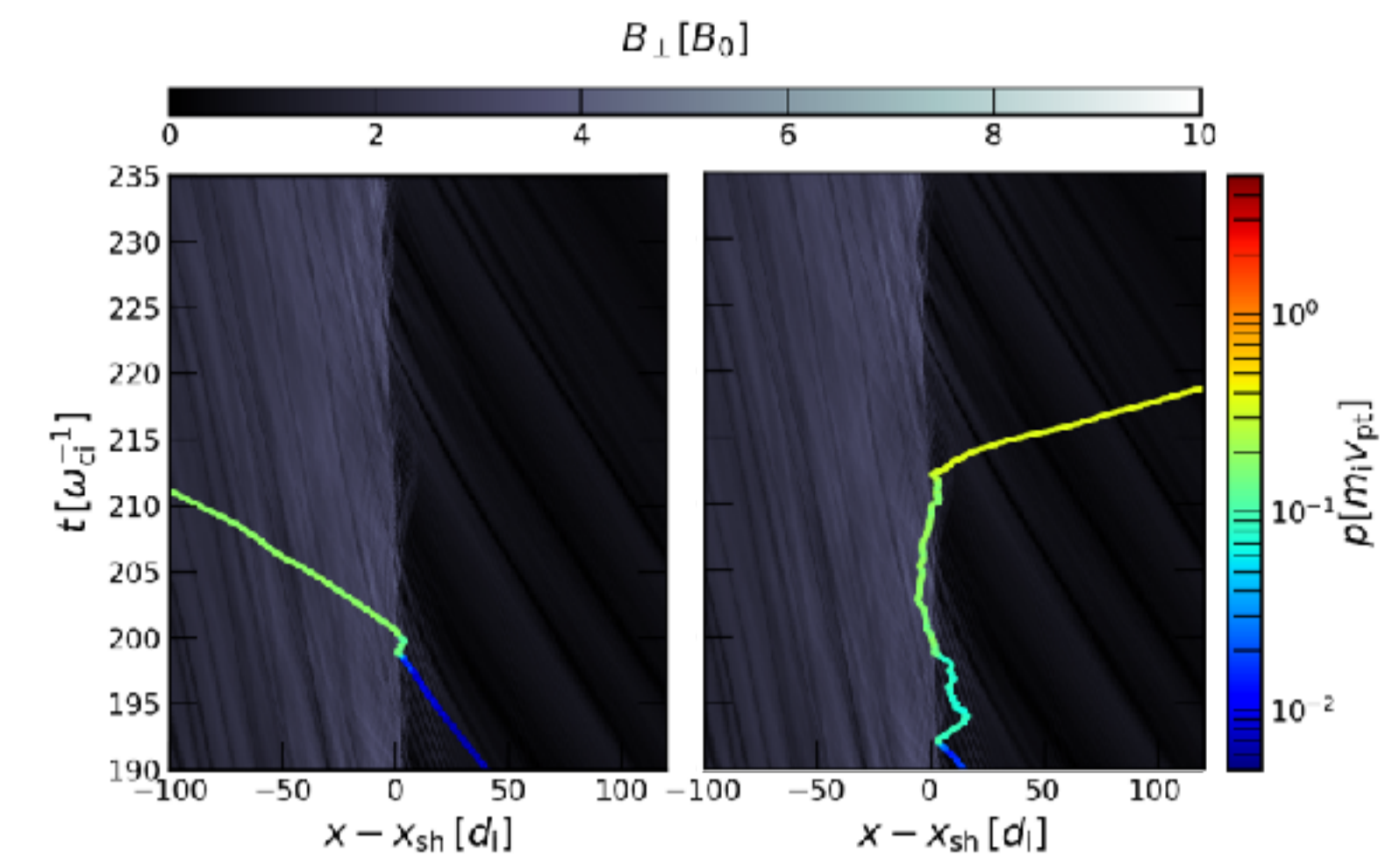


Figure 3. Trajectory of two tracer electrons in  $M_A = 5$  shock.

High Mach number shock

- Particles are confined near the shock.
- Maximum energy is increasing.

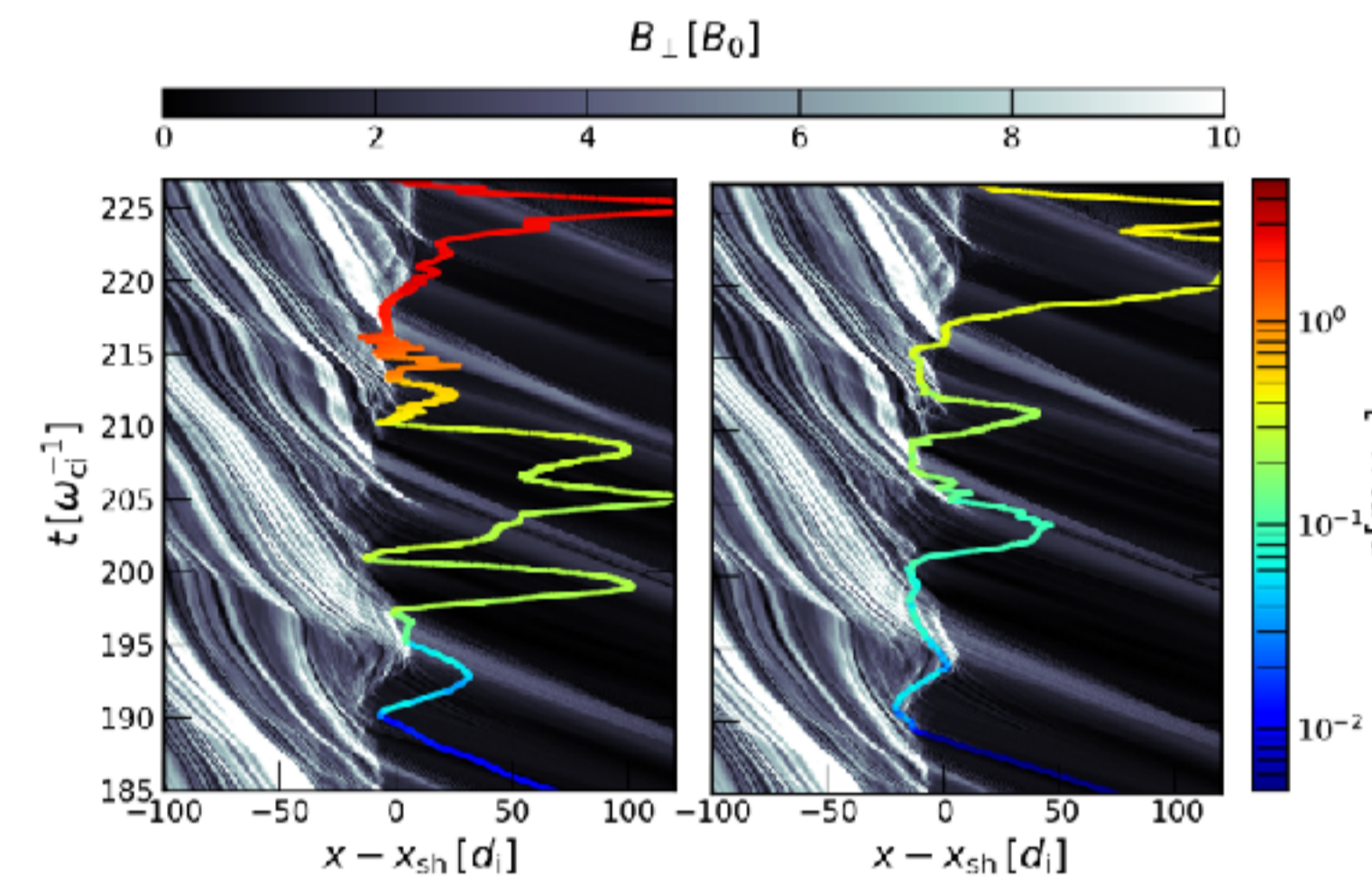


Figure 4. Trajectory of two tracer electrons in  $M_A = 20$  shock.

## Shock structure

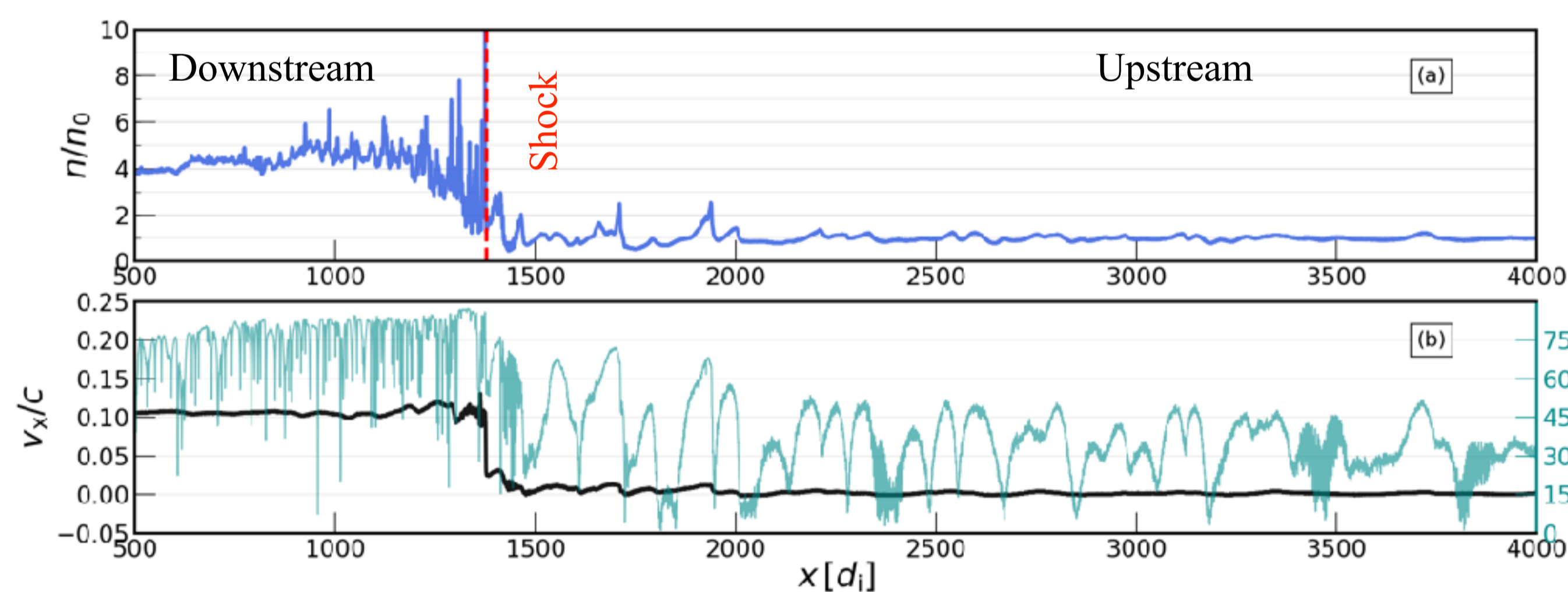


Figure 1. Profiles of density, the  $x$ -component of the plasma speed, and the effective shock inclination  $\theta_{Bn} = \cos^{-1}(B_x/|\mathbf{B}|)$  at  $t = 275 \omega_{ci}^{-1}$ . Run parameters:  $v_{pt}/c = 0.1$ ,  $M_A = 20$ ,  $M_s = 40$ , and  $m_i/m_e = 100$ .

Our investigation focuses on quasi-parallel shocks, where the inclination of the upstream  $\mathbf{B}$  field relative to shock normal is initialized with  $\theta_{Bn} = 30^\circ$ . Shock is launched using a piston moving at a speed  $v_{pt}$  in the upstream frame.

- Plasma (thermal + cosmic rays) density profile is similar to MHD shocks.
- Interactions between thermal and energetic particles produce instabilities.
- Self-generated electromagnetic turbulence re-arranges  $\theta_{Bn}$ .

## Dependence on Mach number

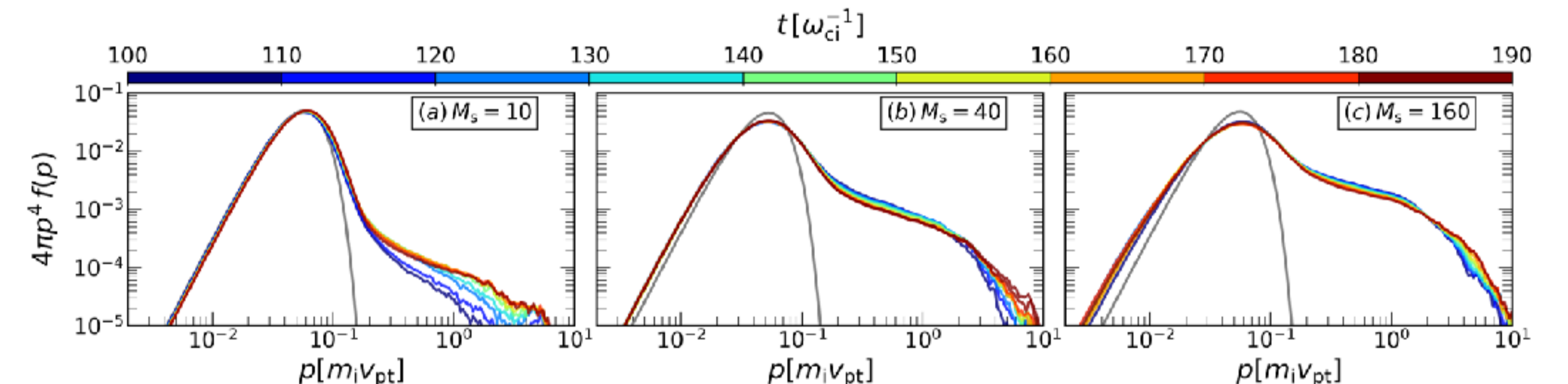


Figure 5. Downstream electron spectra for  $M_s = 10, 40$ , and  $160$  shocks (electron sonic Mach numbers 1, 4, and 16 respectively)

- Acceleration works for both subsonic and supersonic electrons.
- Nonthermal fraction reduces with smaller  $M_s$  shocks.
- Low  $M_A$  shocks show a cut-off at smaller momentum  $p/(m_i v_{pt})$ .
- High Mach number shocks accelerate electrons to highest energies.

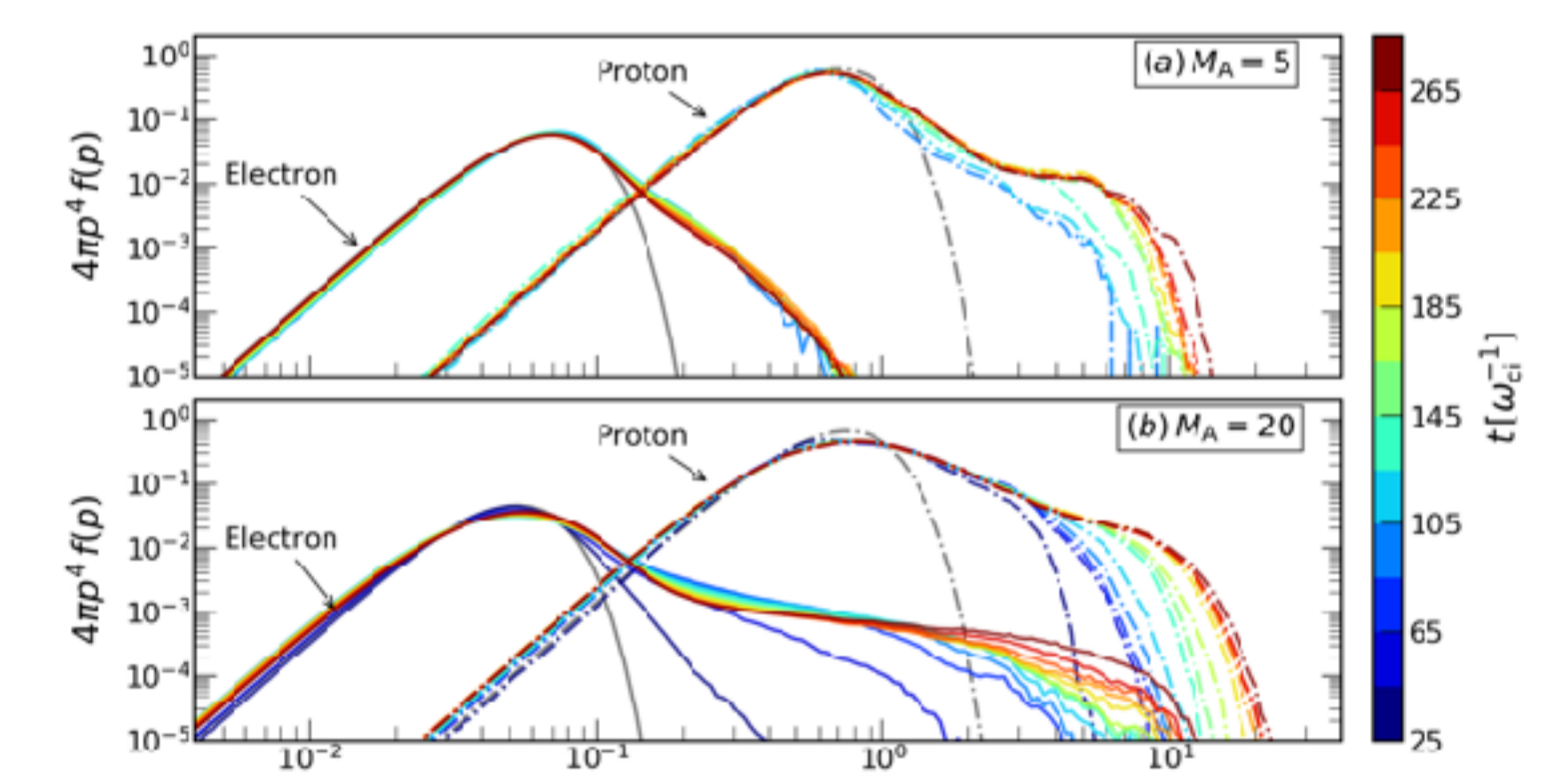


Figure 6. Comparison of downstream electron and proton spectra between  $M_A = 5$  (panel a) and  $M_A = 20$  (panel b) shocks.

## Momentum distribution and spectra

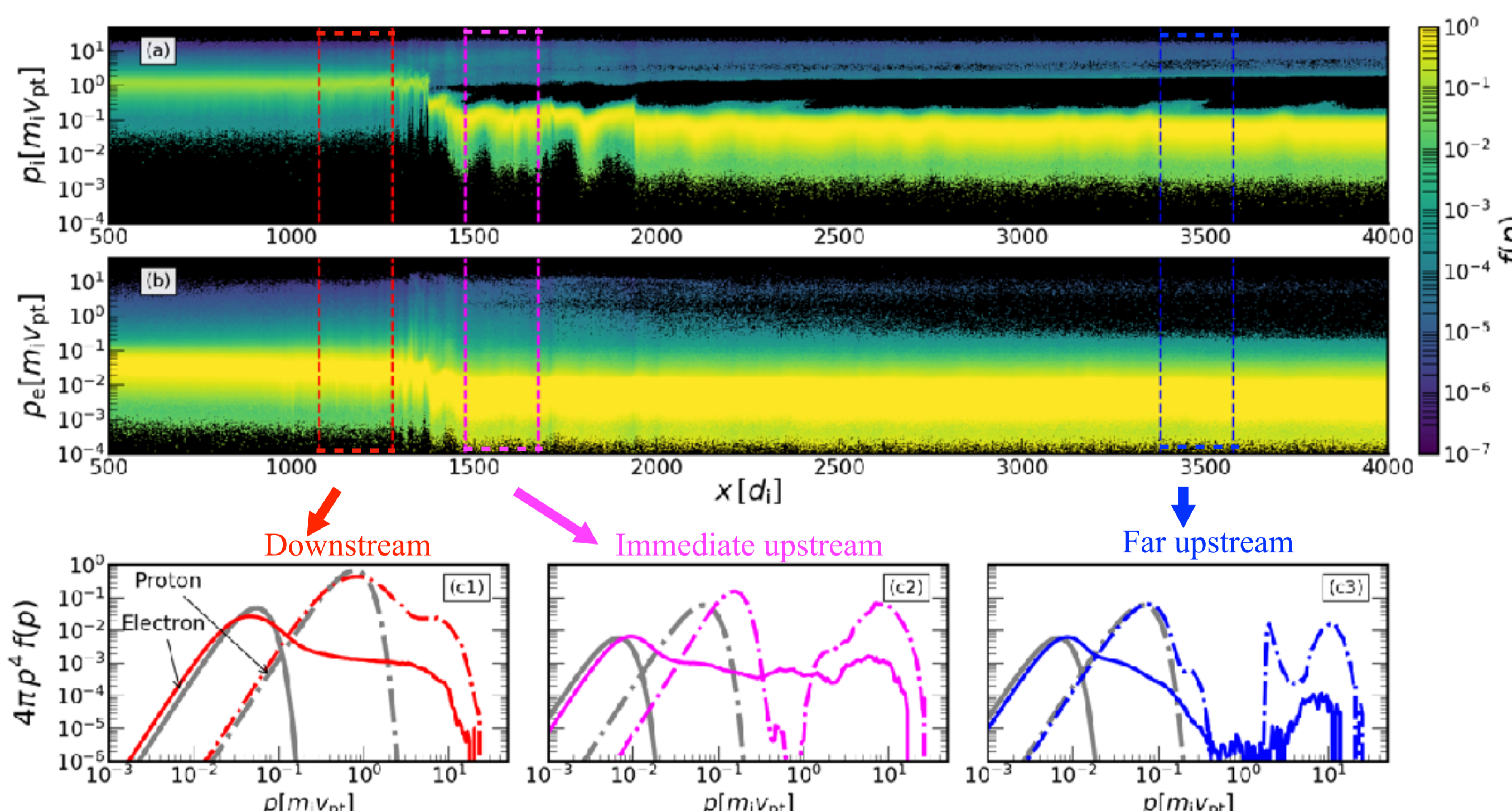


Figure 2. The  $x - |\mathbf{p}|$  diagram at  $t = 275 \omega_{ci}^{-1}$ . Panels (a) and (b) stand for proton and electron, respectively. Panels (c1)-(c3) represent the spectra of protons (dash-dotted line) and electrons (solid line).

- Upstream proton distributions contain thermal and nonthermal populations.
- The electron distribution is made of three populations: thermal, current compensating super-thermal, and nonthermal electrons.

## Summary

- Non-relativistic quasi-parallel shocks can produce nonthermal electrons efficiently (e.g., Figs 1 and 2).
- Reflectivity of a shock and the electromagnetic fluctuations both are crucial for electron DSA injection (Figs 3 and 4).
- Subsonic electrons can also participate in the DSA (Fig 5).
- Due to lack of large-amplitude modes in low  $M_A$  shocks, the spectra show cut-off at smaller energies (Figs 3 and 6).
- In high Mach number shocks, the nonthermal tail keeps growing to higher energies (Figs 4 and 6).

## References including:

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